Infinitesimal Transformations In A Projective Symmetric Finsler Space

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Abstract: In this paper we find infinitesimal projective and special projective transformations with certain theorems and relations in a projective symmetric Finsler space.

Keywords: Infinitesimal transformations, Projective symmetric Finsler space MSC2010:53B40

I. INTRODUCTION

Takano, K. [8] studied the projective motion in a Riemannian space with bi-recurrent curvature. Sinha, R.S [6] discussed the infinitesimal projective transformation in a Finsler space, Yano, K. and Nagano, T. [10] have defined projective conformal transformation in a Riemannian space. Pande, H.D. and Kumar, A [4] have also discussed special infinitesimal projective transformation in a Finsler space and have obtained certain results.

In the present paper studies have been carried out with references to infinitesimal projective and special projective transformations and accordingly certain theorems and relations have been obtained in a projective symmetric Finsler space.

In view of the Berwald's covariant derivative the Lie derivatives of a tensor field $T_j^i(x, \dot{x})$ and the connection parameter $G_{jk}^i(x, \dot{x})$ are given as under:

(1.1)
$$\begin{array}{l} \pounds_{v} T_{j}^{i}(x,\dot{x}) \underline{\underline{def}} T_{j(h)}^{i} v^{h} + \left(\dot{\partial}_{n} T_{j}^{i} \right) v_{(s)}^{h} \dot{x}^{s} T_{j}^{h} v_{(h)}^{i} + T_{h}^{i} v_{(j)}^{h} \\ \text{And} \end{array}$$

(1.2)
$$\mathbf{\pounds}_{v} \mathbf{G}_{jk}^{i}(\mathbf{x}, \dot{\mathbf{x}}) = \mathbf{v}_{(j)(k)}^{i} \mathbf{H}_{jkh}^{i} \mathbf{v}^{h} + \mathbf{G}_{sjk}^{i} \mathbf{v}_{(r)}^{s} \dot{\mathbf{x}}^{r}$$

where $\mathbf{H}_{jkh}^{i}(\mathbf{x}, \dot{\mathbf{x}})$ is defined by [I-(14.6)].

We also have the following commutation formulae

(1.3)
$$\dot{\partial}_{l} \left(\pounds_{v} T_{j}^{i} \right) - \pounds_{v} \left(\pounds_{v} T_{j}^{i} \right) = 0,$$
(1.4)
$$\pounds_{v} T_{j(k)}^{i} - \left(\pounds_{v} T_{j}^{i} \right)_{(k)} = T_{j}^{h} \pounds_{v} G_{kh}^{i} - T_{h}^{i} \pounds_{v} G_{kj}^{h} - \left(\partial_{h} T_{j}^{i} \right) \pounds_{v} G_{ks}^{h} \dot{x}^{s}$$
And

(1.5)
$$(\mathbf{\pounds}_{v}\mathbf{G}_{jh}^{i})_{(k)} - (\mathbf{\pounds}_{v}\mathbf{G}_{kh}^{i})_{(j)} = \mathbf{\pounds}_{v}\mathbf{Q}_{hjk}^{i} + (\mathbf{\pounds}_{v}\mathbf{G}_{kl}^{r})\mathbf{G}_{rjh}^{i}\dot{\mathbf{x}}^{l} - (\mathbf{\pounds}_{v}\mathbf{G}_{ll}^{r})\dot{\mathbf{x}}^{l}\mathbf{G}_{rhk}^{i}\dot{\mathbf{x}}^{l}.$$

In view of the projective covariant derivative the Lie derivative of $T_j^i(x, \dot{x})$ and the connection parameter $\Pi_{jk}^i(x, \dot{x})$ are given by: (1.6) $\pounds_v T_j^i(x, \dot{x}) = T_{j((r))}^i v^r + (\dot{\partial}_s T_j^i) v_{((r))}^s \dot{x}^r - T_j^r v_{((r))}^i + T_r^i v_{((j))}^r$

(1.6) $\pounds_{v} T_{j}^{l}(x, \dot{x}) = T_{j((r))}^{l} v^{r} + (\partial_{s} T_{j}^{l}) v_{((r))}^{s} \dot{x}^{r} - T_{j}^{r} v_{((r))}^{l} + T_{r}^{l} v_{((j))}^{r}$ And

(1.7)
$$\pounds_{v}\Pi^{i}_{mk} = v^{i}_{((m))((k))} + Q^{i}_{mkr}v^{r} + (\dot{\partial}_{r}\Pi^{i}_{mk})v^{r}_{((s))}\dot{x}^{s}$$

The connection parameter Π_{mk}^{1} satisfies the identity

(1.8)
$$\Pi^1_{\text{rkh}} \dot{\mathbf{x}}^{\,\text{h}} = 0.$$

Between the operators \mathbf{f}_v , $\hat{\partial}$ and ((k)), we can obtain the following commutation formulae

(1.9)
$$\hat{\partial}_{1} \left(\pounds_{v} \mathbf{T}_{j}^{i} \right) - \pounds_{v} \left(\hat{\partial}_{1} \mathbf{T}_{j}^{i} \right) = 0$$
(1.10)
$$\left(\pounds_{v} \mathbf{T}_{j}^{i} \right)_{((\mathbf{r}))} - \pounds_{v} \mathbf{T}_{j((\mathbf{r}))}^{i} = \mathbf{T}_{j}^{l} \pounds_{v} \Pi_{\mathbf{r}l}^{i} - \mathbf{T}_{l}^{i} \pounds_{v} \Pi_{\mathbf{r}j}^{l} - \left(\hat{\partial}_{1} \mathbf{T}_{j}^{i} \right) \pounds_{v} \Pi_{\mathbf{r}m}^{l} \dot{\mathbf{x}}^{m}$$

And

(1.11)
$$\left(\pounds_{v} \Pi_{jh}^{i} \right)_{((k))} - \left(\pounds_{v} \Pi_{kh}^{i} \right)_{((j))} = \pounds_{v} Q_{hjk}^{i} + \left(\pounds_{v} \Pi_{kl}^{r} \right) \Pi_{rjh}^{i} \dot{x} - \left(\pounds_{v} \Pi_{jl}^{r} \right) \dot{x}^{1} \Pi_{rhk}^{i}$$

II. SYMMETRIC FINSLER SPACE AND AFFINE MOTION

A symmetric Finsler space is characterised by the vanishing of the covariant derivative of Berwald's curvature tensor field H^i_{hik} .

DEFINITION (2.1)

When the extended point transformation

(2.1)
$$\xi^{x'} = \xi'^{x} = f^{x}(\xi^{v})$$

does not change the fundamental function $L(\xi, \dot{\xi})$ of a Finsler space, that is, when we have

(2.2)
$$L('x, '\dot{\xi}) = L(\xi, \dot{\xi}),$$

we call this extended point transformation a motion in a Finsler space $F_{n}. \label{eq:Finsler}$

DEFINITION (2.2)

When the following transformations $\xi'^{x} = f^{r}(\xi^{v})$

(2.3) and

(2.4)
$$\dot{\xi}^{x} = \left(\partial_{\lambda} f^{x}\right) \dot{\xi}^{\lambda}$$

transform every geodesic into a geodestic and the affine parameter on it into an affine parameter on the deformed geodesic, we call the transformation an affine motion in general affine space of geodesics.

III. NON-AFFINE INFINITESIMAL PROJECTIVE TRANSFORMATION

Let us consider an infinitesimal point transformation

(3.1)
$$\overline{\mathbf{x}^{i}} = \mathbf{x}^{i} + \mathbf{v}^{i}(\mathbf{x}) d\mathbf{t}$$

where $v^i(x)$ determines a non-zero contravariant vector field defined over the domain of the space under consideration and dt is an infinitesimal constant. If the infinitesimal transformation (3.1) transform the system of geodesics of paths into the same system then such a transformation is termed as infinitesimal projective transformation in F_n .

The necessary and sufficient condition in order that an infinitesimal transformation (3.1) be an infinitesimal projective transformation is given by the following equation [3]:

(3.2)
$$\pounds_{v} G^{i}_{jk} = \overline{G}^{i}_{jk} - G^{i}_{jk} = \delta^{i}_{j} p_{j} - g_{jk} g^{il} d_{l}$$

here $p_k(x, \overline{x})$ and $d_1(x, \overline{x})$ are vectors and satisfy the

following identities

(3.3) (a)
$$\dot{\partial}_{j} \mathbf{p} = \mathbf{p}\mathbf{j}$$
, (b) $\mathbf{p}_{hk} = \dot{\partial}_{n}\dot{\partial}_{k}\mathbf{p}$,
(c) $\mathbf{p}_{hk}\dot{\mathbf{x}}^{h} = \mathbf{p}_{k}$, (d) $\mathbf{p}_{hk}\dot{\mathbf{x}}^{h}\dot{\mathbf{x}}^{k} = \mathbf{p}$,
(e) $\dot{\partial}_{j}\mathbf{d} = \mathbf{d}\mathbf{j}$ (f) $\mathbf{d}_{hk} = \dot{\partial}_{h}\dot{\partial}_{k}\mathbf{d}$,
(g) $\mathbf{d}_{hk}\dot{\mathbf{x}}^{h} = \mathbf{d}_{k}$, (h) $\mathbf{d}_{hk}\dot{\mathbf{x}}^{h}\dot{\mathbf{x}}^{k} = \mathbf{d}$

With the help of the commutation formula (1.5), the Liederivative of H^{i}_{hik} is given by

$$(3.4) \quad \pounds_{\mathbf{v}} \mathbf{H}_{hjk}^{i} = \left(\pounds_{\mathbf{v}} \mathbf{G}_{jh}^{i} \right)_{(k)} - \left(\pounds_{\mathbf{v}} \mathbf{G}_{kh}^{i} \right)_{(j)} - \left(\pounds_{\mathbf{v}} \mathbf{G}_{kl}^{r} \right) \mathbf{G}_{rjh}^{i} \dot{\mathbf{x}}^{1} + \left(\pounds_{\mathbf{v}} \mathbf{G}_{jl}^{r} \right) \dot{\mathbf{x}}_{l} \mathbf{G}_{rhk}^{i}$$

With the help of the equation of Lie-derivative, the equation (3.4) assumes the form

$$(3.5) \quad \pounds_{v} H_{hjk}^{i} = \delta_{j}^{i} p_{h(k)} - \delta_{k}^{i} p_{h(j)} + \delta_{h}^{i} p_{j(k)} = \delta_{h}^{i} p_{k(j)} - g_{jh} g^{il} d_{l(k)} + g_{kh} g^{il} d_{l(j)} + g_{kl} g^{rm} G_{rjh}^{i} d_{m} \dot{x}^{l} - g_{jl} g^{rm} G_{rkh}^{i} d_{m} \dot{x}^{l}.$$

Multiplying (3.5) by $\dot{x}^{h}\dot{x}^{j}$ and noting the equation (5.2) and the homogeneity property of $H^{i}_{hik}(x,\dot{x})$, we get

3.6)
$$\mathbf{\pounds}_{v} \mathbf{H}_{k}^{i} = 2\dot{x}^{i} p_{(k)} - \delta_{k}^{i} p_{(j)} \dot{x}^{j} - \dot{x}^{i} p_{k(j)} \dot{x}^{j} - g_{jh} g^{il} d_{l(k)} \dot{x}^{h} \dot{x}^{j} + g_{kh} g^{il} d_{l(j)} \dot{x}^{h} \dot{x}^{j}.$$

Contracting (3.6) with respect to the indices i and k we get

(3.7)
$$\pounds_{v} H = -p_{(j)} \dot{x}^{j} + \frac{1}{n-1} \left\{ d_{j} \dot{x}^{j} - g_{hk} g^{il} de_{(i)} \dot{x}^{h} \dot{x}^{j} \right\}$$

With the help of equation (3.6) and (3.7), we get

$$(3.8) \quad \pounds_{v} H_{k}^{i} - \pounds_{v} H \delta_{k}^{i} = 3\dot{x}^{i} p_{(x)} - \delta_{k}^{i} p_{(j)} \dot{x}^{j} - \dot{x}^{i} p_{k(j)} \dot{x}^{j} + g_{kh} g^{il} d_{l(j)} \dot{x}^{h} \dot{x}^{j} - \frac{1}{n-1} \left\{ d_{(k)} \dot{x}^{i} + (2-n) g_{jh} g^{il} d_{l(k)} \dot{x}^{h} \dot{x}^{j} \right\}$$

Differentiating (3.8) partially with respect to \dot{x}^r and other contracting the resulting equation with respect to the indices i & r, we get,

$$(3.9) \left(\pounds_{v} \dot{\partial}_{r} H_{k}^{r} - \pounds_{v} \dot{\partial}_{k} H \right) = (3n+2)p_{(k)} - (n+3)p_{k(j)} + d_{k(j)} \dot{x}^{j} + g_{kh} g^{rl} \dot{x}^{h} \left\{ d_{rl(j)} \dot{x}^{j} + d_{l(r)} \right\} - \frac{5-n}{n-1} d_{l(k)} + 2 \dot{x}^{h} \dot{x}^{j} \frac{C_{sr}^{l}}{g_{rs}} \left\{ \frac{2-n}{n-1} g_{rh} d_{l(k)} - g_{kh} d_{l(j)} \right\}$$

In view of the equation (3.8) and (3.9), the Lie derivative of $W_i^i(x, \dot{x})$ can be written as:

$$(3.10) \ \pounds_{v} W_{k}^{i} = \frac{1}{n+1} \left[p_{(k)} \dot{x}^{i} + 2p_{k(j)} \dot{x}^{i} \dot{x}^{j} + \frac{4-n}{n-1} d_{(k)} \dot{x}^{i} - \dot{x}^{i} \left\{ d_{k(j)} \dot{x}^{j} + g_{kh} g^{rl} \dot{x}^{h} \left(d_{rl(j)} \dot{x}^{j} + d_{l(r)} \right) + \right. \right]$$

(

$$+2\dot{x}^{h}\dot{x}^{j}\frac{C_{sr}^{l}}{g_{rs}}\left(\frac{2-n}{n-1}g_{rh}d_{l(k)}-g_{kh}d_{l(j)}\right)\right\}$$

$$-\delta_{k}^{i}p_{(j)}\dot{x}^{j}+g_{kh}g^{il}d_{l(j)}\dot{x}^{h}\dot{x}^{j}+\frac{2-n}{n-1}g_{jh}g^{il}d_{l(k)}\dot{x}^{h}\dot{x}^{j}.$$

Applying the commutation formula to the projective deviation tensor $W_i^1(x, \dot{x})$, we get

$$(3.11) \mathfrak{L}_{v} W_{j(r)}^{i} = \left(\mathfrak{L}_{v} W_{j}^{i} \right)_{(r)} = W_{j}^{h} \mathfrak{L}_{v} G_{rh}^{i} - W_{h}^{i} \mathfrak{L}_{v} G_{jr}^{h} - \left(\dot{\partial}_{h} W_{j}^{i} \right) \mathfrak{L}_{v} G_{rs}^{h} i^{s}.$$

Substituting the value of $\pounds_{v} G_{ik}^{1}$ from (3.2) in (3.11) and using the equations (3.3), we get

$$(3.12) \ \pounds_{v} W_{j(r)}^{i} - \left(\pounds_{v} W_{j}^{i}\right)_{(r)}^{\prime} = W_{j}^{h} \left(\delta_{r}^{i} p_{h} - g_{rh} g^{il} dl\right) - W_{r}^{i} p_{j} - 2W_{j}^{i} p_{r} + g^{hl} d_{l} \left\{W_{h}^{i} g_{jr} + \left(\dot{\partial}_{h} W_{j}^{i}\right) g_{rs} \dot{x}^{s}\right\} - \left(\dot{\partial}_{r} W_{j}^{i}\right) p$$

Contracting (3.12) with respect to the indices i and r and noting we get

Transvecting (3.12) by \dot{x}^r and noting we get

$$(3.14) \left\{ \pounds_{v} W_{j(r)}^{i} - \left(\pounds_{v} W_{j}^{i} \right)_{(r)} \right\} \dot{x}^{r} = W_{j}^{h} \dot{x}^{i} p_{h} - 4W_{j}^{i} p - W_{j}^{h} g_{rh} g^{il} d_{l} \dot{x}^{r} + g^{hl} d_{l} \dot{x}^{r} \left\{ W_{h}^{i} g_{gr} + \left(\dot{\partial}_{h} W_{j}^{i} \right) g_{rs} \dot{x}^{s} \right\}$$

At this stage, if we now assume that the Finsler space F_n symmetric one i.e., $W_{i}^{i}(r) = 0$? then under this is assumption $\pounds_v W^i_{j(r)} = 0$ will always hold and therefore equations (3.13) and (3.14) will respectively assume the following form :

(3.15)
$$\begin{pmatrix} \pounds_{v} W_{j}^{i} \end{pmatrix}_{(r)} \dot{x}^{r} = (2-n) W_{j}^{h} p_{h} + W_{j}^{h} d_{h} - g^{hl} d_{l} - g^{hl} d_{l} \left\{ W_{h}^{i} g_{jr} + (\dot{\partial}_{h} W_{j}^{i}) g_{rs} \dot{x}^{s} \right\}$$
And

$$(3.16) \quad \left(\pounds_{v}W_{j}^{i}\right)_{(r)}\dot{x}^{r} = -W_{j}^{h}\dot{x}^{i}p_{h} + W_{j}^{i}p + W_{j}^{h}g_{rh}g^{il}d_{l}\dot{x}^{r} - g^{hl}d_{l}\dot{x}^{r} \left\{W_{h}^{i}g_{jr} + \left(\dot{\partial}_{h}W_{j}^{i}\right)g_{rs}\dot{x}^{s}\right\}$$

With the help of equation (3.15) and (3.16) we eliminate the term $W_i^h p_h$ and the result of elimination gives

$$(3.17) \quad \mathbf{M}_{j}^{i} = \dot{\mathbf{x}}^{i} \left[\mathbf{W}_{j}^{h} \mathbf{d}_{h} - \mathbf{g}^{hl} \mathbf{d}_{l} \left\{ \mathbf{W}_{h}^{r} \mathbf{g}_{jr} + \left(\dot{\partial}_{h} \mathbf{W}_{j}^{r} \right) \mathbf{g}_{rs} \dot{\mathbf{x}}^{s} \right\} \right] + 2(2-n) \left[4 \mathbf{W}_{j}^{i} \mathbf{p} + \mathbf{W}_{j}^{h} \mathbf{g}_{rh} \mathbf{g}^{il} \mathbf{d}_{l} \dot{\mathbf{x}}^{r} - \mathbf{g}^{hl} \mathbf{d}_{l} \dot{\mathbf{x}}^{r} \left\{ \mathbf{W}_{h}^{i} \mathbf{g}_{jr} + \left(\dot{\partial}_{h} \mathbf{W}_{j}^{i} \right) \mathbf{g}_{rs} \dot{\mathbf{x}}^{s} \right\} \right]$$

Where

(3.18)
$$\mathbf{M}_{j}^{i} \underline{\underline{\det}} \left(\mathbf{\pounds}_{v} \mathbf{W}_{j}^{i} \right)_{(r)} \dot{\mathbf{x}}^{i} + (2 - n) \left(\mathbf{\pounds}_{v} \mathbf{W}_{j}^{i} \right)_{(r)} \dot{\mathbf{x}}^{r}.$$

If the Finsler F_n admits a projective affine motion, then the equation

$$(3.19) \qquad \qquad \mathbf{\pounds}_{v} = \mathbf{G}_{jk}^{i} = \mathbf{0}$$

must hold.

Therefore, with the help of equation (3.2) and (3.19), it is clear that the vectors $p(x, \dot{x})$ and $d(x, \dot{x})$ must separately vanish. Thus, we can state

THEOREM (3.1)

If an F_n admits a non-affine infinitesimal projective transformation such that the Berwarld's covariant derivatives of W_i^1 remains invariant then the equation (3.17) holds.

THEOREM (3.2)

If an F_{n} admits an affine infinitesimal projective transformation such that the covariant derivative of W_i^1 remains invariant then $M_i^1 = 0$.

THEOREM (3.3)

If a symmetric Finsler space F_n admits a non-affine infinitesimal projective transformation then the equation (3.17) necessarily holds.

THEOREM (3.4)

In a symmetric Finsler space, if an infinitesimal projective transformation is an affine one then the equation $M_{i}^{1} = 0$ necessarily holds.

IV. INFINITESIMAL SPECIAL PROJECTIVE **TRANSFORMATION DEFINITION (4.1)**

The necessary and sufficient condition in order that the infinitesimal point transformation (3.1) be an infinitesimal special projective transformation is given by

(4.1)
$$\pounds_{\mathbf{v}}\Pi^{i}_{jk} = \overline{\Pi}^{i}_{jk} - \Pi^{i}_{jk} = \delta^{i}_{j}b_{k} + \delta^{i}_{j}b_{k} - g_{gk}g^{il}c_{l}$$

Here $b_k(x, \dot{x})$ and $C_1(x, \dot{x})$ are vector which satisfy the following relations

(4.2) (a)
$$\dot{\partial}_{j}b = b_{j}$$
, (b) $b_{hk} = \dot{\partial}_{h}\dot{\partial}_{k}b$,
(c) $b_{hk}\dot{x}^{h} = b_{k}$, (d) $b_{hk}\dot{x}^{h}\dot{x}^{k} = b$,
(e) $\dot{\partial}_{j}c = c_{j}$, (f) $c_{jk} = \dot{\partial}_{j}\dot{\partial}_{k}c$,
(g) $c_{hk}\dot{x}^{h} = c_{k}$ (h) $c_{nk}\dot{x}^{h}\dot{x}^{k} = c$.

With the help of the equation (4.1), (4.2) and the commutation formula (1.11), the Lie-derivative of the projective entity Q_{hjk}^1 is given by

$$(4.3) \quad \pounds_{v} Q_{hjk}^{i} = \delta_{j}^{i} b_{h((k))} + \delta_{h}^{i} b_{j((k))} - g_{jh} g^{il} c_{l((kk))} - g_{jh} g^{il} c_{l((kk))} - g_{jh} g^{il} c_{l} - g_{jh} g^{il}_{((k))} c_{l} - \delta_{k}^{i} b_{h((j))} - \delta_{h}^{i} b_{k((j))} + g_{kh} g^{il} c_{l((j))} + g_{kh} (g^{il} c_{l} + g^{il} c_{l((j))} + g_{kh} (g^{il} c_{l} + g^{il} c_{l((j))} + g_{kh} (g^{il} c_{l} + g^{il} c_{l}) + g_{kh} (g^{il} c_{l} + g^$$

$$\begin{split} &+g_{kh}g^{il}_{((j))}c_e - \delta^r_k b\Pi^i_{rjh} + g_{kl}g^{rm}c_m\Pi^i_{rjh}\dot{x}^i + \\ &+\delta^r_j b\Pi^i_{rhk} - g_{jl}g^{rm}c_m\dot{x}^l\Pi^i_{rhk}. \end{split}$$

Multiplying (4.3) by $\dot{x}^h \dot{x}^j$ and using the equation (1.8) and the fact that $Q^i_{hjk} \dot{x}^h = Q^i_{jk}$, we get

$$\begin{array}{ll} (4.4) \quad \pounds_{v} Q_{k}^{i} = 2\dot{x}b_{((k))} - \delta_{k}^{i}b_{((j))}\dot{x}^{j} + \dot{x}^{h}\dot{x}^{j} & \left[g_{kh}\left\{g_{((j))}^{il}c_{l} + g^{il}c_{l((j))}c_{l} + g^{il}c_{l((j))}\right\} - g_{gh}\left\{g_{((k))}^{il}c_{l} + g^{il}c_{l((k))}\right\} - g_{jh((k))}g^{il}c_{l}\right] \end{array}$$

Contracting (4.3) with respect to the indices i and k and then multiplying the resulting equation thus obtaining by

 $\dot{x}^{h}\dot{x}^{j}$ and using (1.8), we get (4.5)

$$\pounds_{v} Q_{hj} \dot{x}^{h} \dot{x}^{j} = (1-n) b_{((j))} \dot{x}^{j} + c_{((j))} \dot{x}^{j} + g^{il} c_{l} \dot{x}^{h} \dot{x}^{j} \left\{ g_{il((j))} - g_{l} \right\}$$

$$-g_{jh}\dot{x}^{h}\dot{x}^{j}\left\{g^{il}c_{l((i))} + g^{il}_{((i))}c_{l}\right\} + g_{ih}g^{il}_{((j))}c_{l}\dot{x}^{h}\dot{x}^{j}.$$

We now eliminate the term $b_{((j))}\dot{x}^{j}$ from the equation (4.4) and (4.5) and get

Where

(4.7)
$$\mathbf{M}_{k}^{i}(\mathbf{x}, \dot{\mathbf{x}}) \underline{\underline{\det}}(i-n) \pounds_{v} \mathbf{Q}_{k}^{i} + \delta_{k}^{i} \pounds_{v} \mathbf{Q}_{hj} \dot{\mathbf{x}}^{h} \dot{\mathbf{x}}^{j}$$

Applying the commutation formula (1.9) to the projective derivation tensor $W_j^i(x, \dot{x})$ and using the equation (4.2), we get

$$\begin{aligned} & (4.8) \quad \left(\pounds_{\mathbf{v}} \mathbf{W}_{j}^{i} \right)_{((\mathbf{r}))} - \pounds_{\mathbf{v}} \mathbf{W}_{j((\mathbf{r}))}^{i} = \mathbf{W}_{j}^{l} \delta_{\mathbf{r}}^{i} \mathbf{b}_{l} - \mathbf{W}_{j}^{l} \mathbf{g}_{\mathbf{r}l} \mathbf{g}^{ip} \mathbf{c}_{p} - \mathbf{W}_{\mathbf{r}}^{i} \mathbf{b}_{j} + \\ & + \mathbf{W}_{l}^{i} \mathbf{g}_{\mathbf{r}j} \mathbf{g}^{lp} \mathbf{c}_{p} - \left(\dot{\partial}_{\mathbf{r}} \mathbf{W}_{j}^{i} \right) \mathbf{b} - 2 \mathbf{W}_{j}^{\mathbf{r}} \mathbf{b}_{\mathbf{r}} - \\ & - \left(\dot{\partial}_{\mathbf{r}} \mathbf{W}_{j}^{i} \right) \mathbf{g}_{m} \mathbf{g}^{lp} \mathbf{c}_{p} \dot{\mathbf{x}}^{m}. \end{aligned}$$

Transvecting (4.8) by \dot{x}^r we get

10)
$$\left\{ \left(\pounds_{v} W_{j}^{i} \right)_{((r))} - \pounds_{v} W_{j((r))}^{i} \right\} \dot{x}^{r} = W_{j}^{l} b_{l} \dot{x}^{i} - - 4 W_{j}^{i} b - W_{j}^{l} g_{rl} g^{lp} c_{p} \dot{x}^{r} + + W_{l}^{i} g_{rj} g^{lp} c_{p} \dot{x}^{r} - \left(\dot{\partial}_{l} W_{j}^{i} \right) g_{rm} g^{lp} c_{p} \dot{x}^{r} \dot{x}^{m}$$

Let us now suppose that the infinitesimal special projective transformation (4.1) leaves in variant the projective covariant derivative of the projective deviation tensor i.e.

(4.11)
$$\pounds_{v} W_{j((r))}^{i} = 0.$$

In view of (4.11), equations (4.9) and (4.10) can be alternatively expressed in the following forms

(4.12)
$$\left(\pounds_{v} W_{j}^{i} \right)_{((j))} = (n-2) W_{j}^{l} b_{l} - W_{j}^{l} c_{l} + g^{lp} c_{p} \left\{ W_{l}^{i} g_{rj} - (\dot{\partial}_{l} W_{j}^{i}) g_{lm} \dot{x}^{m} \right\}$$

And

(4.

$$\begin{array}{c} (4.13) & \left(f_{v} W_{j}^{i} \right)_{((r))} \dot{x}^{r} = W_{j}^{l} b_{l} \dot{x}^{i} - 4 W_{j}^{i} b - W_{j}^{l} g_{rl} g^{ip} c_{p} \dot{x}^{r} + \\ & + W_{l}^{i} g_{rj} g^{lp} c_{p} \dot{x}^{r} - \\ & - \left(\partial_{l} W_{j}^{i} \right) g_{m} g^{ip} c_{p} \dot{x}^{r} \dot{x}^{m} \end{array}$$

respectively.

Eliminating the term $W_j^l b_l$ from equations (4.12) and (4.13), we get

4.14)
$$\begin{split} B_{j}^{i}(x,\dot{x}) &= x^{i} \left[-W_{j}^{l}c_{l} + g^{lp}c_{p} \left\{ W_{l}^{i}g_{kj} - \left(\dot{\partial}_{l}W_{j}^{i}\right)g_{im}\dot{x}^{m} \right\} \right] - \\ &- (n-2) \left[-4W_{j}^{i}b - W_{j}^{i}g_{rl}g^{ip}c_{p}\dot{x}^{r} + W_{l}^{i}g_{rj}g^{lp}c_{p}\dot{x}^{r} - \\ &- \left(\partial_{i}W_{j}^{i}\right)g_{rm}g^{lp}c_{p}\dot{x}^{r}\dot{x}^{m} \right] \end{split}$$

Where

(4.15)
$$\mathbf{B}_{j}^{i}(\mathbf{x}, \dot{\mathbf{x}}) \underline{\det} \left(\pounds_{\mathbf{v}} \mathbf{W}_{j}^{i} \right)_{((i))} \dot{\mathbf{x}}^{i} - (n-2) \left(\pounds_{\mathbf{v}} \mathbf{W}_{j}^{i} \right)_{((nr))} \dot{\mathbf{x}}^{r}.$$

If the Finsler space F_n admits special projective affine motion, then the equation

(4.16)
$$\pounds_{v} \Pi^{1}_{jh} = 0 \text{ must hold.}$$

Hence equations (4.1) and (4.16) enables us to infer that the vectors $b(x, \dot{x})$ and $c(x, \dot{x})$ must separately vanish, Therefore, we can state

THEOREM (4.1)

If an F_n admits a non-affine infinitesimal special projective transformation such that the projective covariant derivative of W_j^i remains invariant then equation (4.15) holds.

THEOREM (4.2)

If an F_n admits an affine infinitesimal special projective transformation such that the covariant derivative of W_j^i remains invariant then B_j^i given in (4.15) should vanish.

In particular, if we now suppose that the Finsler space F_n in the symmetric one i.e. $W_{j((r))}^i = 0$ holds then equation (4.11) always holds. Therefore, we can state

THEOREM (4.3)

If a symmetric Finsler space F_n admits a non-affine infinitesimal special projective transformation then the equation (4.14) holds necessarily.

THEOREM (4.4)

In a symmetric Finsler space $F_n \mbox{ if an infinitesimal special projective transformation is an affine one then the equation$

 $B_{i}^{i} = 0$ necessarily holds.

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